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LETTER TO THE EDITOR

A new integrable modified Korteweg–de Vries equation with one half degree of nonlinearity

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Abstract. An alternative generalized equation which incorporates the modified Korteweg–de Vries (MKdV) and KdV equations is proposed and from it a new integrable system of the MKdV type is found.

Integrable systems play a fundamental role in nonlinear science. Since the discovery of the complete integrability of the KdV equation [1], a special interest has been growing in the exactly solvable systems. In this letter, we report a new integrable nonlinear evolution equation which has the form

$$16u_t + u_{xxx} + 30u^{1/2}u_x = 0. \tag{1}$$

This equation is a special case of the following generalized MKdV equation

$$n^2u_t + u_{xxx} + (n+1)(n+2)u^{2/n}u_x = 0 \tag{2}$$

where n is a positive integer. For $n=4$, (2) gives (1). Clearly, (2) also incorporates the MKdV equation ($n=1$) and the KdV equation ($n=2$). However, unlike the generalized KdV equation [2]

$$u_t + u_{xxx} + u^n u_x = 0 \tag{3}$$

which takes the KdV equation as the basic equation, (2) takes the MKdV equation as the base.

In the following, we investigate the integrability of (2) through the Painlevé method [3]. The substitution of $u = q^n$ leads to the equation

$$n^2q^2q_t + (n-1)(n-2)q^3 + 3(n-1)qq_xq_{xx} + q^2q_{xxx} + (n+1)(n+2)q^4q_x = 0. \tag{4}$$

We assume that

$$q = \phi^\alpha \sum_{j=0}^{\infty} V_j \phi^j \tag{5}$$

where $V_j = V_j(x, t)$, ϕ is a singular manifold. Substituting (5) into (2), it is found that

$$\alpha = -1 \quad V_0^2 = -\phi_x^2 \tag{6}$$

and the recursion relation takes the form

$$(j+1)[j-(n+2)][j-(2n+2)]V_0^2\phi^3 V_j = F(V_{j-1}, \dots, V_0; \phi_t, \phi_x, \dots). \tag{7}$$

Since the detailed form of F is very complicated, it will be presented elsewhere [4]. It is clear from (7) that the resonances occur at

$$j = -1 \quad n + 2 \quad 2n + 2. \quad (8)$$

From (8) we find that

$$j = 0: V_0^2 = -1 \quad (9)$$

$$j = 1: V_1 = 0 \quad (10)$$

$$j = 2: V_2 = V_0 \psi_t / 6 \quad (11)$$

$$j = 3: (n - 1) V_3 = 0 \quad (12)$$

$$j = 4: (n - 1)(n - 2)(10 V_4 + 7 V_0 V_2^2) = 0 \quad (13)$$

$$j = 5: 6(n - 3)(2n - 3) V_5 + n^2 V_{2t} = 0 \quad (14)$$

$$j = 6: (n - 2)(n - 4)(14 V_6 + 31 V_2^3 / 5) = 0 \quad (15)$$

$$j = 7: 10(2n - 5)(n - 5) V_7 + (76n^2 - 327n + 302) V_0 V_2 V_5 = 0 \quad (16)$$

$$j = 8: 18(n - 3)(n - 6) V_8 + n^2 V_{5t} + (191n^2 + 141n - 2762) 3 V_0 V_2^4 / 100 \\ + 6(4n^2 - 15n + 2) V_0 V_2 V_6 = 0 \quad (17)$$

$$j = 10: -22(n - 4)(n - 8) V_{10} + n^2(2 V_0 V_2 V_5 - V_7)_t + (-58n^2 + 246n - 206) V_0 V_5^2 \\ + (-32n^2 + 114n + 236) V_0 V_2 V_8 + (-641n^2 + 3117n + 2078) 3 V_2^5 / 100 \\ + 4(-8n^2 + 60n - 67) V_2^2 V_6 = 0 \quad (18)$$

where we have used the reduced ansatz [5]

$$\phi = x + \psi(t) \quad \phi_x = 1 \quad V_j = V_j(t).$$

When $n = 1$ and $n = 2$, the compatibility conditions at $j = n + 2$ and $j = 2n + 2$ are satisfied identically and the $m\kappa\lambda\nu$ and $\kappa\lambda\nu$ equations possess the Painlevé property. This is a well known fact [3]. However, the result that (2) with $n = 4$, i.e. (1), also possesses the Painlevé property, and is thus identified as integrable, is new. This is easy to see from (15) and (18) since the resonances occur at $j = 6$ and $j = 10$ respectively. The integrability of (2) with general n is difficult to investigate through the Painlevé method since (7) becomes very complicated for large n .

References

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